

Similarly.

$$\frac{\rho_B v_{1B} r_B - \rho_D v_{1D} r_D}{\Delta r} \rightarrow \frac{\partial \rho v_{1r}}{\partial r}$$

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$$\frac{\rho_A v_{2A} - \rho_C v_{2C}}{\Delta r} \rightarrow \frac{\partial \rho v_{2\theta}}{\partial \theta}$$

Total differential continuity equation

$$\frac{\partial \rho}{\partial t} + \underbrace{\frac{1}{r} \frac{\partial \rho v_{1r}}{\partial r} + \frac{1}{r} \frac{\partial \rho v_{2\theta}}{\partial \theta}}_{\text{table book or appendix B:}} = \text{div}(\rho \vec{v})$$

Apparently then

$$\boxed{\frac{\partial \rho}{\partial t} + \text{div}(\rho \vec{v}) = 0}$$

is differential continuity equation

Method 2

Rigorous, and quicker: convert surface integrals into volume ones, using divergence theorem

divergence theorem: If \mathcal{V} is a volume and A is the complete terminating surface of \mathcal{V} , and \vec{v} any vector:

$$\boxed{\int_{\mathcal{V}} \text{div} \vec{v} \cdot d\mathcal{V} = \oint_A \vec{v} \cdot \vec{n} \, dA}$$

Index notation

$$\int_{\mathcal{V}} \frac{\partial v_i}{\partial x_i} d\mathcal{V} = \oint_A v_i n_i dA$$

But summation is not needed:

$$\boxed{\int_{\mathcal{V}} \frac{\partial f}{\partial x_i} d\mathcal{V} = \oint_A f n_i dA}$$

i.e. going from surface integrals to volume integrals: *

$$n_i \rightarrow \frac{\partial}{\partial x_i} \text{ [entire rest of the integrand]}$$

$$dA \rightarrow d\mathcal{V}$$

Apply to continuity:

$$\frac{dM_{cv}}{dt} + \int_{A_{cv}} \rho \vec{v} \cdot \vec{n} dA = 0$$

or

$$\frac{d}{dt} \int_{V_{cv}} \rho dV + \int_{A_{cv}} \rho v_i n_i dA = 0$$

div theorem

$$\int_{V_{cv}} \frac{\partial \rho v_i}{\partial x_i} dV = 0$$

change order of integration

$$\int_{V_{cv}} \left[\frac{\partial \rho}{\partial t} + \frac{\partial \rho v_i}{\partial x_i} \right] dV = 0$$

this can only be true for any control volume V_{cv} if the integrand is everywhere 0

$\frac{\partial \rho}{\partial t} + \frac{\partial \rho v_i}{\partial x_i} = 0$ differential continuity in conservation form
 $\underbrace{\quad}_{\text{div } \rho \vec{v}}$ in index notation

Apply to momentum

$$\frac{dL_{cv}}{dt} + \int \rho \vec{v} \cdot \vec{n} \vec{v} dA = \int \underline{\underline{\tau}}_{vis}^T \vec{n} dA - \int p \vec{n} dA + \int \rho \vec{g} dV$$

make sure to use unique index names!

$$\frac{d}{dt} \int \rho v_i dV + \int \rho v_j n_j v_i dA = \int \tau_{ji}^v n_j dA - \int p n_i dA + \int \rho g_i dV$$

$$\int \frac{\partial \rho v_i}{\partial t} dV + \int \frac{\partial \rho v_j v_i}{\partial x_j} dV = \int \frac{\partial \tau_{ji}^v}{\partial x_j} dV - \int \frac{\partial p}{\partial x_i} dV + \int \rho g_i dV$$

Force per unit volume

$\frac{\partial \rho v_i}{\partial t} + \frac{\partial \rho v_j v_i}{\partial x_j} = \frac{\partial \tau_{ji}^{vis}}{\partial x_j} - \frac{\partial p}{\partial x_i} + \rho g_i$

differential momentum in conservation form

instead: $\frac{\partial \tau_{ji}^{vis}}{\partial x_j} = 0$

Apply to energy:

Let internal energy per unit mass be e
 Kinetic energy per unit mass is $\frac{1}{2}V^2 = \frac{1}{2}|\vec{V}|^2$

Treat gravity as a force doing work, rather than an energy per unit mass and height.

$$E_{CV} = \int_{\text{CV}} \rho (e + \frac{1}{2}V^2) dV$$

$$\frac{dE_{CV}}{dt} + \int_{A_{CV}} \rho \vec{v} \cdot \vec{n} (e + \frac{1}{2}v^2) dA = \int_{A_{CV}} \rho \vec{v} \cdot \vec{n} dA + \int_{A_{CV}} \vec{v} \cdot \vec{\tau} dA + \int_{A_{CV}} \rho \vec{v} \cdot \vec{S} dA$$

work done by pressure forces

↓
work done by gravity on fluid particles

↓
work done by pressure and viscous forces on surface A_{CV} per unit time

$$\frac{\partial \rho (e + \frac{1}{2}v_i v_i)}{\partial t} + \frac{\partial \rho v_j (e + \frac{1}{2}v_i v_i)}{\partial x_j} = - \frac{\partial p v_j}{\partial n_j} + \frac{\partial \tau_{ij}^{visc} v_i}{\partial x_j} + \rho v_i g_i$$